### <span id="page-0-0"></span>The Duality of Trace and Determinant

Noah Bergam

August 15th, 2023

Noah, CC '25 [Columbia UMS Summer 2023](#page-25-0) August 15th, 2023 1/26

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# **Outline**







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### <span id="page-2-0"></span>**Tensors**

The tensor product  $V \otimes W$  consists of expressions of the form:

$$
v_1\otimes w_1+v_2\otimes w_2+...+v_n\otimes w_n
$$

This ⊗ is multi-linear, i.e. linear in each component.

$$
(\lambda_1v_1+v_2)\otimes w_1=\lambda_1(v_1\otimes w_1)+(v_2\otimes w_1) \qquad \qquad (1)
$$

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# <span id="page-3-0"></span>Contextualizing Tensors

- Cartesian Product  $V \times W$ : tuples of the underlying set.
- Direct Product  $V \times W$ : tuples as a vector space, coordinate-wise.
- Direct Sum  $V \oplus W$ : tuples with finitely many nonzero components.
- **Tensor Product**  $V \otimes W$ : tuples with multi-linearity
- Wedge Product  $V \wedge V$ : tuples with multi-linearity and anti-symmetry.

(Note: the difference between direct sum and direct product only emerges for infinite sums/products.)

(Note: you cannot wedge different vector spaces, since this would make anti-commutativity ill-defined.)

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# <span id="page-4-0"></span>Exterior Power

### Definition (Exterior Power)

For a vector space V, the n-th exterior power of V, denoted  $\Lambda^n V$ , is spanned by elements of the following form for  $v_1, ..., v_n \in V$ .

 $v_1 \wedge ... \wedge v_n$ 

which obey the multi-linearity and anti-symmetry. For example:

$$
(\lambda v_1 + w) \wedge ... \wedge v_n = \lambda (v_1 \wedge ... \wedge v_n) + (w \wedge ... \wedge v_n)
$$

$$
v_1 \wedge v_2 \wedge ... \wedge v_n = -(v_2 \wedge v_1 \wedge ... \wedge v_n)
$$

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### <span id="page-5-0"></span>Notes on Exterior Product

#### Important Properties:

- $\bigcirc$  (vanishing)  $v \wedge v = 0$
- 2 (associativity)  $(v \wedge w) \wedge x = v \wedge (w \wedge x)$
- **3** (untangling)  $v_{σ(1)}$   $\land$  ...  $\land$   $v_{σ(n)}$  = sgn( $σ$ )( $v_1$   $\land$  ...  $\land$   $v_n$ )

**Warning:** Not every element in  $\Lambda^n V$  is *reducible* to a single wedge product (in general, it is a linear combination of such elements).

# <span id="page-6-0"></span>The k-linear extension

### Definition

For  $A\in \operatorname{\mathsf{End}}(V)$ , the k-linear extension  $\Lambda^N A^k:\Lambda^N V\mapsto \Lambda^N V$  defined as:

$$
\Lambda^m A^k \left( \bigwedge_{j=1}^m v_j \right) = \sum_{s} \bigwedge_{j=1}^m A^{s_j} v_j
$$
\nwhere

\n
$$
s \in \{0, 1\}^n \quad \sum_{j} s_j = 1
$$

 $\Lambda^N A^N$  means we apply  $A$  to each entry (one-dimensional).  $\Lambda^N A^1$  means we apply  $A$  to only one entry (*n*-dimensional).

# <span id="page-7-0"></span>The k-linear extension (examples)

For  $\Lambda^N A^k$  I am applying  $A$  to  $k$  entries of the wedge product.

$$
\Lambda^3 A^1(v_1 \wedge v_2 \wedge v_3) = Av_1 \wedge v_2 \wedge v_3 + v_1 \wedge Av_2 \wedge v_3 + v_1 \wedge v_2 \wedge Av_3
$$

If V is N-dimensional, then dim $(\Lambda^N A^k) = {N \choose k}$ .

$$
\Lambda^3 A^2 (v_1 \wedge v_2 \wedge v_3) = A v_1 \wedge A v_2 \wedge v_3
$$
  
+  $A v_1 \wedge v_2 \wedge A v_3$   
+  $v_1 \wedge A v_2 \wedge A v_3$ 

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# <span id="page-8-0"></span>Determinant and Trace

### Definition

The **determinant** of  $A \in End(V)$  is the number by which any nonzero tensor  $\omega \in \Lambda^N V$  is multiplied when  $\Lambda^N A^N : \Lambda^N V \mapsto \Lambda^N V$  acts on it.

$$
(\Lambda^N A^N)\omega = (\det A)\omega
$$

### **Definition**

The trace of  $A \in End(V)$  is the number by which any nonzero tensor  $\omega \in \Lambda^N V$  is multiplied when  $\Lambda^N A^1 : \Lambda^N \mapsto \Lambda^N$  acts on it.

$$
(\Lambda^N A^1)\omega = (\text{tr } A)\omega
$$

Note: For dim(V) = N,  $\Lambda^N V$  is one-dimensional.  $\Lambda^1 V$  is *n*-dimensional.

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### <span id="page-9-0"></span>Illustration of wedge-based determinant

$$
\Lambda^n A^n \omega = \Lambda^n A^n (v_1 \wedge \ldots \wedge v_n) = (Av_1 \wedge \ldots \wedge Av_n)
$$
  
\n
$$
= \Big( \sum_{j_1=1}^n A_{j_1,1} v_{j_1} \wedge \ldots \wedge \sum_{j_1=1}^n A_{j_n,n} v_{j_n} \Big)
$$
  
\n
$$
= \sum_{j_1=1}^n \ldots \sum_{j_n=1}^n \Big( A_{j_1,1} v_{j_1} \wedge \ldots \wedge A_{j_n,n} v_{j_n} \Big)
$$
  
\n
$$
= \sum_{j_1=1}^n \ldots \sum_{j_n=1}^n (A_{j_1,1} \cdots A_{j_n,n}) \Big( v_{j_1} \wedge \ldots \wedge v_{j_n} \Big)
$$
  
\n
$$
= \sum_{\sigma \in S_n} sgn(\sigma) \prod_{j=1}^n A_{j,\sigma(j)} (v_1 \wedge \ldots \wedge v_n) = det(A) \omega \quad \Box
$$

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### <span id="page-10-0"></span>Illustration of wedge-based trace

$$
\Lambda^n A \omega = \Lambda^n A (v_1 \wedge ... \wedge v_n) = \sum_{i=1}^n v_1 \wedge ... \wedge (A v_i) \wedge ... \wedge v_n
$$

$$
=\sum_{i=1}^n\sum_{j_1=1}^n A_{j_i,i}(v_1\wedge...\wedge v_{j_i}\wedge...\wedge v_n)
$$

This is zero unless  $v_i = i$ . This eliminates the second sum and recovers the usual formula.

$$
=\sum_{i=1}^n A_{ii}(v_1\wedge...\wedge v_n)=\mathrm{tr}(A)\omega
$$

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### <span id="page-11-0"></span>Theorems

### Theorem (Liouville's Formula)

Let  $A \in$  End(V).

$$
\det(\exp(A))=\exp(tr(A))
$$

where  $\exp(A) = 1 + A + \frac{1}{2!}A^2 + \frac{1}{3!}A^2 + ...$  denotes the matrix exponential. More generally,  $det(exp(t \cdot A(t))) = exp(t \cdot tr(A(t)))$ , where t is a formal variable and  $A(t)$  is an operator-valued formal power series.

### Theorem (Jacobi's Formula)

For F(t) an operator-valued formal power series such that  $F^{-1}(t)$  exists:

$$
\partial_t \det(F(t)) = (\det F(t)) \ tr(F^{-1}(t) \cdot \partial_t F(t))
$$

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# <span id="page-12-0"></span>Context and Game Plan

An operator-valued formal power series is just a function  $F(t) = 1 + F_1 t + F_2 t^2 + ...$  where the coefficients  $F_i \in End(V)$ . This is more flexible than just talking about linear operators.

The idea is to represent both  $det(exp(A(t)))$  and  $exp(tr(A(t)))$  as a (formal) power series in  $t$  satisfying some differential equation.

- First we establish some theory on how to solve differential equations for formal power series.
- Then we will guess a suitable differential equation that will enable us to prove the identity.

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# <span id="page-13-0"></span>Characterization of  $exp(tA)$

#### Lemma

The operator-valued function  $F(t) = \exp(tA)$  is the unique solution to the following differential equation.

$$
\partial_t F(t) = F(t)A \qquad F(0) = 1_V
$$

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# <span id="page-14-0"></span>Proof of Characterization of  $exp(tA)$

#### Lemma

For  $A \in End(V)$ . the operator-valued function  $F(t) = \exp(tA)$  is the unique solution to the following differential equation.

$$
\partial_t F(t) = F(t)A \qquad F(0) = 1_V
$$

#### Proof.

Since  $F(0)=1$ , we know  $F(t)=1+F_1t+F_2t^2+...$ Note  $F'(0) = A = F_1 A$ ,  $F''(0) = A^2 = 2F_2$ ,  $F'''(0) = A^3 = 6F_2$ , etc. Matching coefficients, we find:  $F(t) = 1 + At + \frac{1}{2!}A^2t^2 + \frac{1}{3!}A^3t^2 + ... = \exp(tA).$ 

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# <span id="page-15-0"></span>Leibniz Rule for Power Series

#### Lemma

If  $\phi(t)$  and  $\psi(t)$  are power series in t with coefficients from  $\Lambda^m V$  and  $\Lambda^n V$ respectively, then the Leibniz rule holds, i.e.

$$
\partial_t(\phi \wedge \psi) = (\partial_t \phi) \wedge \psi + \phi \wedge (\partial_t \phi)
$$

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# <span id="page-16-0"></span>Proof of Leibniz Rule for Power Series

#### Lemma

If  $\phi(t)$  and  $\psi(t)$  are power series in t with coefficients from  $\Lambda^m V$  and  $\Lambda^n V$ respectively, then the Leibniz rule holds, i.e.

$$
\partial_t(\phi \wedge \psi) = (\partial_t \phi) \wedge \psi + \phi \wedge (\partial_t \phi)
$$

### Proof.

Due to linearity of derivative and the fact that power series can be differentiated term by term, just check for  $\phi=t^2\omega_1$  and  $\psi=t^b\omega_2.$ 

$$
\partial_t(\phi \wedge \psi) = (a+b)t^{a+b-1}\omega_1 \wedge \omega_2
$$

$$
(\partial_t \phi) \wedge \psi + \phi \wedge (\partial_t \psi) = at^{a-1} \omega_1 \wedge t^b \omega_2 + t^a \omega_1 \wedge bt^{b-1} \omega_2
$$

### <span id="page-17-0"></span>Inverse

#### Lemma

The inverse of a formal power series  $\phi(t)$  exists iff  $\phi(0) \neq 0$ .

### Proof.

If  $\phi(0) \neq 0$  then  $\phi(t) = \phi(0) + t\psi(t)$  with  $\psi$  another power series. Then we can construct the inverse explicitly:

$$
\frac{1}{\phi(t)} = \frac{1}{\phi(0)} \frac{1}{(1 + \frac{t\psi(t)}{\phi(0)})} = \sum_{n=0}^{\infty} (-1)^n \phi(0)^{-n-1} (t\psi(t))^n
$$

This is because  $1 = (1 + x)(1 - x + x^2 - x^3 + ...)$  for formal x.

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### Jacobi

### Lemma (Jacobi's Formula)

If  $A(t)$  is an invertible operator-valued formal power series:

$$
\partial_t \det(A(t)) = \det(A) tr(A^{-1} \partial_t A)
$$

Note that the determinant and trace of  $A(t)$  still makes sense because  $A(t)$  is still an operator (just expressed as an infinite sum of operators).

Note: this formula can be written in a lot of different ways (Cf. Wikipedia).

 $\leftarrow$   $\Box$ 

# **Jacobi**

### Lemma (Jacobi's Formula)

If A is an invertible operator-valued formal power series:

$$
\partial_t \det(A(t)) = \det(A) tr(A^{-1} \partial_t A)
$$

### Proof.

Apply definition of determinant and the Leibniz rule established earlier.

$$
(\partial_t \det(A(t)))(\omega) = \partial_t (\det(A(t))\omega) = \partial_t (A v_1 \wedge ... \wedge A v_n)
$$

$$
= \sum_{k=1}^n Av_1 \wedge ... \wedge (\partial_t A)v_k \wedge ... \wedge Av_n
$$

(We want to write this as a trace.)

# Jacobi Proof, Continued

### Proof, continued.

Invoke the algebraic complement of  $A(t)$ , given by  $\det(A(t)) \cdot A^{-1}(t)$  for invertible A (there is a general formula as well). Think of it like the adjoint.

$$
\sum_{k=1}^n Av_1 \wedge ... \wedge (\partial_t A)v_k \wedge ... \wedge Av_n = \sum_{k=1}^n v_1 \wedge ... \wedge (\tilde{A}\partial_t Av_k) \wedge ... \wedge v_n
$$

Note that the right-hand side is a trace:  $\Lambda^n(\tilde{A}\partial_t A)^1 ({\sf v}_1\wedge...\wedge {\sf v}_n).$  This gives us the desired identity.

$$
\partial_t \det(A) = \text{tr}(\tilde A \partial_t A) = \text{tr}(\det(A) A^{-1} \partial_t A) = \det(A) \text{tr}(A^{-1} \partial_t A)
$$

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# Proof of Liouville

Let  $F(t) = \exp(tA)$ .  $F(0) = 1$  so it is invertible, so we can apply Jacobi's:

$$
\partial_t \det(F(t)) = \det(F(t)) \cdot \text{tr}(F^{-1}\partial_t F)
$$

By the characterization of  $F(t) = \exp(tA)$ , we have  $F^{-1}(\partial_t F) = F^{-1}(FA) = (F^{-1}F)A = A.$ 

$$
\partial_t \det(F(t)) = \det(F(t)) \cdot \text{tr}(A)
$$
  
Let  $f(t) = \det(F(t)) \qquad \partial_t f(t) = f(t) \cdot \text{tr}(A)$ 

By the characterization,  $f(t) = \exp(t \cdot \text{tr}(A))$ . Hence:

$$
\det(\exp(tA))=\exp(t\cdot \text{tr}(A))\quad \Box
$$

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### <span id="page-22-0"></span>Related Identities

**Generalization of Liouville's:** For  $p \le n = \dim(V)$  with  $A \in End(V)$ . Liouville's is the special case where  $p = n$ 

$$
\Lambda^p(\exp(tA))^p=\exp(t(\Lambda^pA^1))
$$

**Sylvester's Theorem:** For  $A: V \mapsto W$ ,  $B: W \mapsto V$ , we have:

$$
\det(I_V + BA) = \det(I_W + AB)
$$

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# <span id="page-23-0"></span>An intuitive taste of Jacobi

A nice vignette in Arnold's ODE textbook. **Observation:** As  $\epsilon \to 0$  we have:

$$
\det(I+\epsilon A)=1+\epsilon\mathrm{tr}(A)+O(\epsilon^2)
$$

We can view this easily via the eigenvalues of A, call them  $\lambda_1, ..., \lambda_n$ .

$$
\det(I+\epsilon A)=\prod_{i=1}^n(1+\epsilon\lambda_i)
$$

Note that the zeroth order term is 1. The second term is (by Vieta)  $\epsilon(\sum_{i=1}^n\lambda_i)=\epsilon\mathrm{tr}(A).$  Rest of the terms are order  $\epsilon^2.$ 

# <span id="page-24-0"></span>Conclusion

There are a number of ways to view the duality of trace and determinant.

- Definitions and constructions.
- Theorems and analytical connections.

Careful understanding of the basic constructions (wedge product, power series, differential equation for exp, etc) was key.

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<span id="page-25-0"></span>

Linear Algebra via Exterior Products. Sergei Winitzki. Section 4.5. Ordinary Differential Equations. VI Arnold. Section 16.

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